

# **Neutrinos and Cosmology**

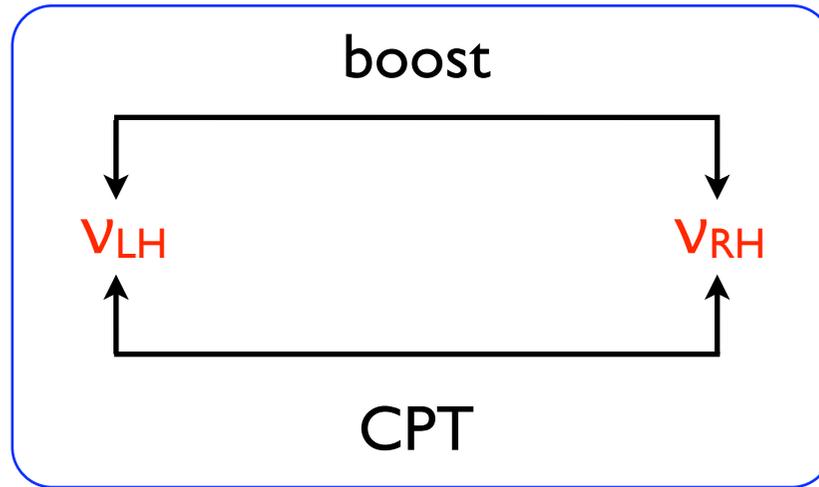
- ❑ ***Open questions in neutrino physics***
- ❑ ***Laboratory neutrino program***
- ❑ ***The inner space/outer space connections***

## Background

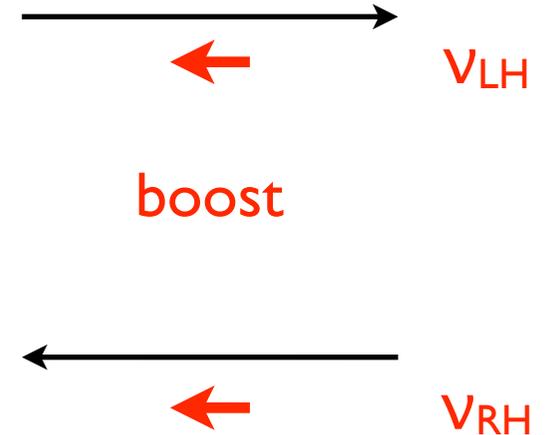
- Neutrino properties one of the four “big Questions” identified by the Astro 2010 Panel on Cosmology and Fundamental Physics
  - the mechanism responsible for inflation
  - the nature of dark energy and/or other explanations for the accelerating universe
  - the nature of the dark matter
  - $\nu$  properties: masses, mixing angles, lepton no. of the cosmos
- New  $\nu$  properties are our one discovery beyond the minimal SM
  - Dirac masses require a RHed  $\nu$  field, absent from the SM
  - Majorana masses correspond to the simplest “effective operator” correction to the SM, requiring a new physics scale  $1/M_{\text{new}}$
- Initial discoveries were made in astrophysics, but now an exciting interface between lab and cosmological  $\nu$  physics has been established
  - there are critical  $\nu$  properties issues, important to the progress of the field, that may only be answerable in cosmology

$\nu$ 's richer mass structure, compared  
other SM fermions

Majorana  
limit:

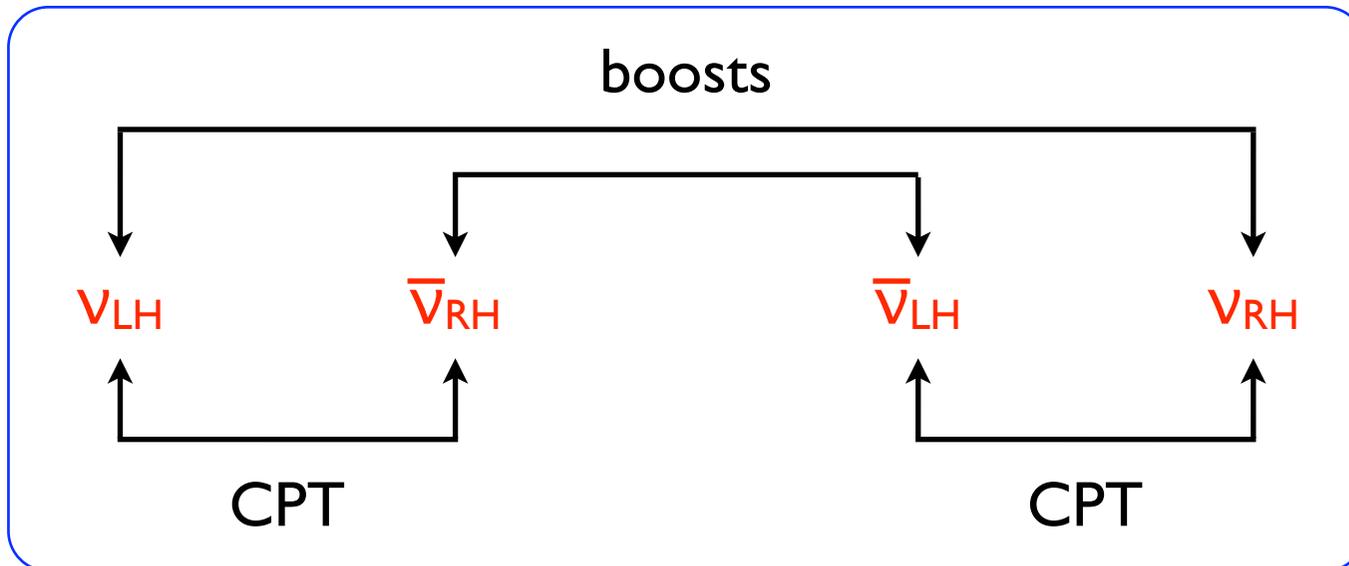


Lorentz invariance



in general, linear  
combinations  
of the two

Dirac  
limit:



The absence of any additively conserved charge allows the addition of

$$L_m(x) \Rightarrow M_D \bar{\Psi}(x) \Psi(x) + (\bar{\Psi}_L^c(x) M_L \Psi_L(x) + \bar{\Psi}_R^c(x) M_R \Psi_R(x) + h.c.)$$

to give the more general matrix

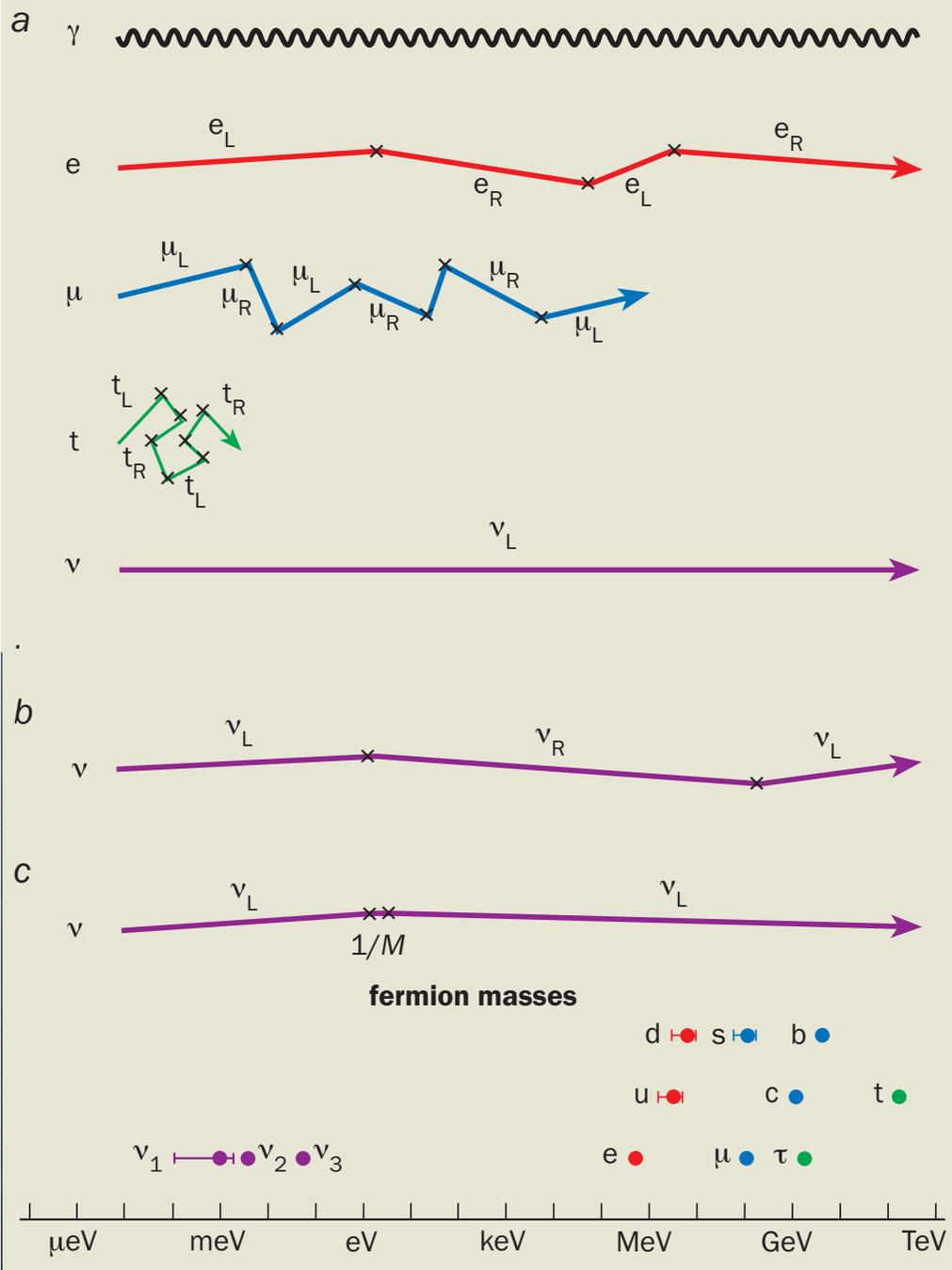
$$(\bar{\Psi}_L^c, \bar{\Psi}_R, \bar{\Psi}_L, \bar{\Psi}_R^c) \begin{pmatrix} 0 & 0 & M_L & M_D^T \\ 0 & 0 & M_D & M_R^\dagger \\ M_L^\dagger & M_D^\dagger & 0 & 0 \\ M_D^* & M_R & 0 & 0 \end{pmatrix} \begin{pmatrix} \Psi_L^c \\ \Psi_R \\ \Psi_L \\ \Psi_R^c \end{pmatrix}$$

which has a number of interesting properties, including

- the introduction of  $M_L, M_R$  breaks the global invariance  $\Psi \rightarrow e^{i\alpha} \Psi$  associated with a conserved lepton number
- while  $M_L$  couples  $\Psi_L$ , weak isospin requires a novel mass mechanism

$$M_L \sim \frac{\langle \phi \rangle^2}{M_{new}}$$

## 2 Neutrinos meet the Higgs boson



Hitoshi Muryama's  $\nu$  mass cartoon

standard model masses

light Dirac neutrino

LHed Majorana neutrino

← the anomalous  $\nu$  mass scale!

This was one reason the initial oscillation results were so exciting:

- give the  $\nu$  an  $M_D$  typical of other SM fermions
- take  $M_L \sim 0$ , in accord with  $\beta\beta$  decay
- assume  $M_R \gg M_D$  as we have not found new RHed physics at low E

$$\begin{pmatrix} 0 & m_D \\ m_D & m_R \end{pmatrix} \Rightarrow m_\nu^{\text{light}} \sim m_D \left( \frac{m_D}{m_R} \right) \quad \text{the needed small parameter!}$$

- take  $m_\nu \sim \sqrt{m_{23}^2} \sim 0.05$  eV and  $m_D \sim m_{\text{top}} \sim 180$  GeV

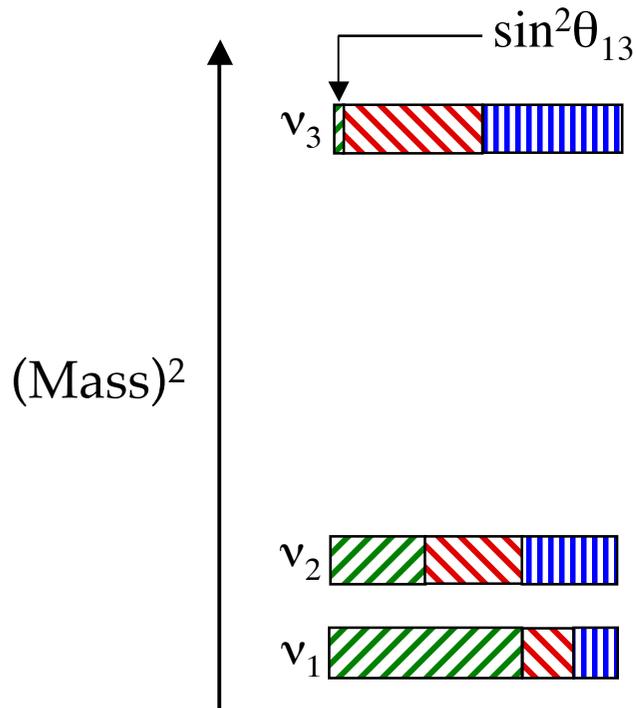
$$\Rightarrow m_R \sim 0.3 \times 10^{15} \text{ GeV}$$

*The deduced  $\nu$  atmospheric mass difference is consistent with a novel mass generation mechanism, not shared by other SM fermions, that the data suggest might be characteristic of the GUT scale*

# The Laboratory/Cosmology Program

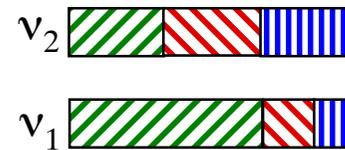
Important questions unanswered

## Hierarchy



Normal

or



Inverted

to within current experimental accuracy, a mixing angle of  $45^\circ$

indications it may be nonzero

  $\nu_e [ |U_{ei}|^2 ]$

  $\nu_\mu [ |U_{\mu i}|^2 ]$

  $\nu_\tau [ |U_{\tau i}|^2 ]$

(matter effects seen only in the solar case)

(artwork: Boris Kayser)

# The mixing

(where we have additional blanks to fill in)

knowns:  $\theta_{12}, \theta_{23}$

known unknowns:  $\theta_{13}, \delta, \phi_1, \phi_2$

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix} \begin{pmatrix} \nu_1 \\ e^{i\phi_1} \nu_2 \\ e^{i\phi_2} \nu_3 \end{pmatrix}$$

$$= \begin{pmatrix} 1 & & \\ & c_{23} & s_{23} \\ & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & s_{13}e^{-i\delta} \\ & 1 \\ -s_{13}e^{i\delta} & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} \\ -s_{12} & c_{12} \\ & & 1 \end{pmatrix} \begin{pmatrix} \nu_1 \\ e^{i\phi_1} \nu_2 \\ e^{i\phi_2} \nu_3 \end{pmatrix}$$

atmospheric

$\nu_e$  disappearance

solar

results:  $\theta_{23} \sim 45^\circ$

$\sin \theta_{13} \leq 0.17$

$\theta_{12} \sim 30^\circ$

$\Delta_{12}$

$|\Delta_{23}|$

$\text{sign}[\Delta_{23}]$

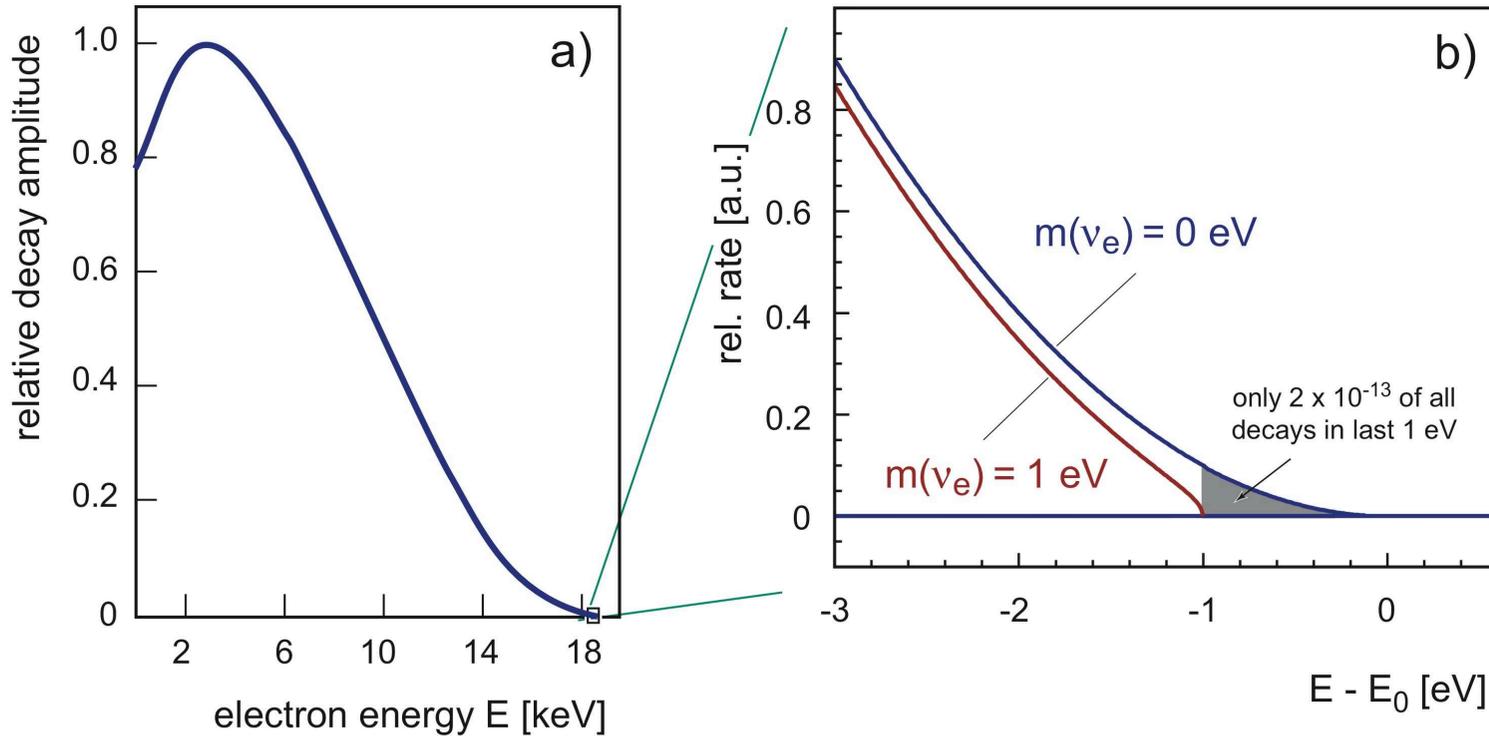
absolute scale

If there are no surprises, progress is needed in four areas

- 1) absolute mass scale
- 2) lepton number and the mass mechanism
- 3) the hierarchy: new matter effects in LBNE
- 4) CP violation and associated questions about the size of  $\theta_{13}$

A 15-20 year lab program is taking shape ~ \$2-3B

## Absolute $\nu$ mass: the one identified component of DM



tritium  $\beta$  decay is running into intrinsic limits due to feasible source intensities and detector resolution

$$\langle m_\nu \rangle_{\text{tritium}} = \sum_i |U_{ei}|^2 m_\nu^2(i)$$

present limit  $\langle m_\nu \rangle_{\text{tritium}} < 2.2 \text{ eV}$

Mainz & Troitzk

KATRIN's goal is to reach 250 meV, with  $5\sigma$  exclusion at 350 meV



the measurement is clean, and one could get lucky ... but cosmology may provide our best hope of reaching the 50 meV level

## lepton number and the mass mechanism: neutrinoless $\beta\beta$ decay

$$\langle m_\nu^{\text{Maj}} \rangle = \sum_{i=1}^{2n} \lambda_i U_{ei}^2 m_i \quad \text{or} \quad \left\langle \frac{1}{m_\nu^{\text{heavy}}} \right\rangle = U_{ei}^2 \frac{1}{m_i^{\text{heavy}}}$$

analogous to the search for the Higgs

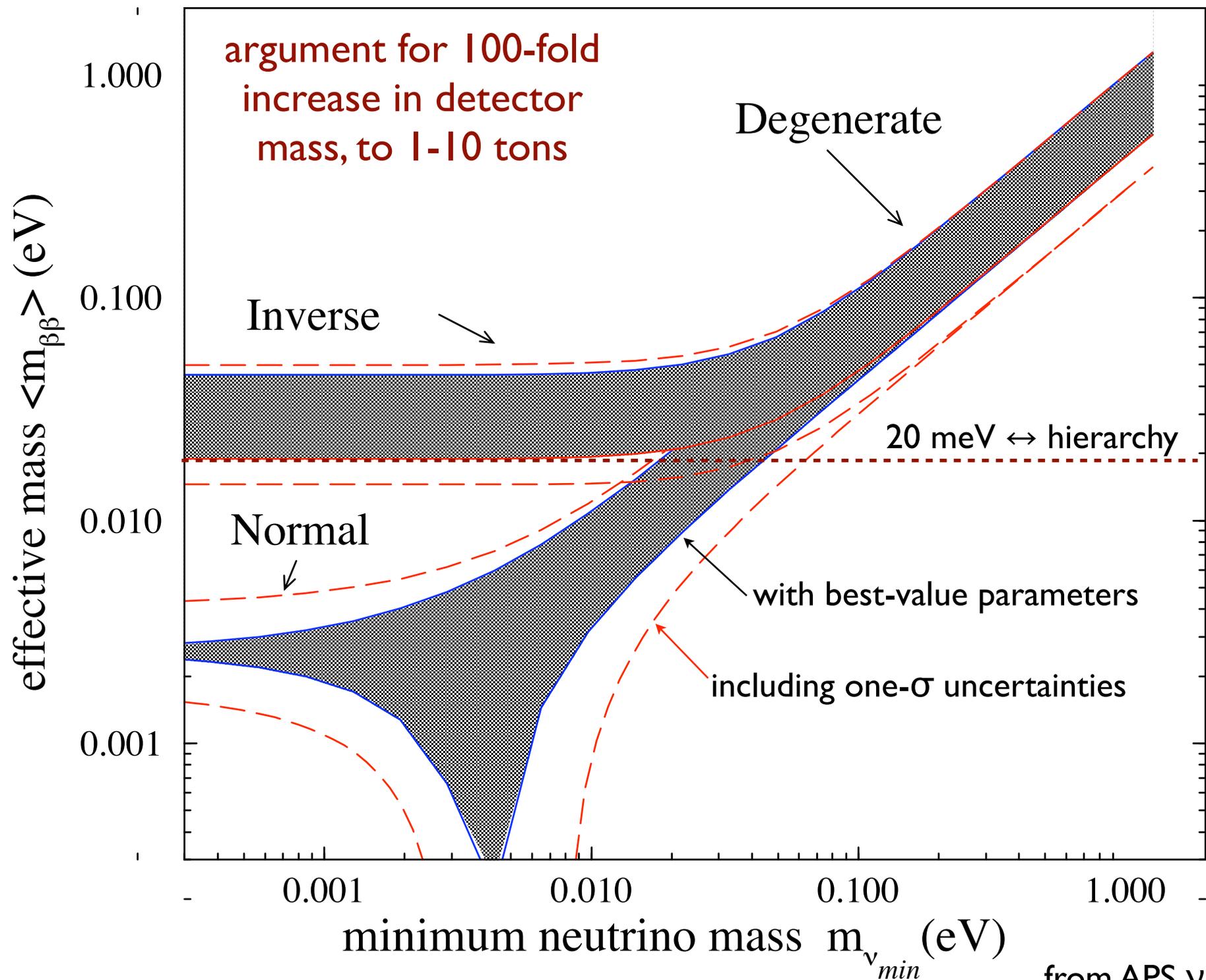
- a mass mechanism connected to the simplest effective SM operator
- indirect sensitivity to near-GUT-scale physics
- direct sensitivity to heavy- $\nu$  super-TeV physics

GERDA ( $^{76}\text{Ge}$ ), CUORE ( $^{128}\text{Te}$ ) currently limit

$$\langle m_\nu^{\text{Maj}} \rangle < (0.3 - 1.0) \text{ eV} \quad \left\langle \frac{1}{m^{\text{heavy}}} \right\rangle < \frac{1}{10^4 \text{ TeV}}$$

**Should be attacked with urgency and at a more elevated scale**

My view: a test of lepton number nonconservation, and possibly of the hierarchy. More ambiguity in deducing  $\nu$  mass scale



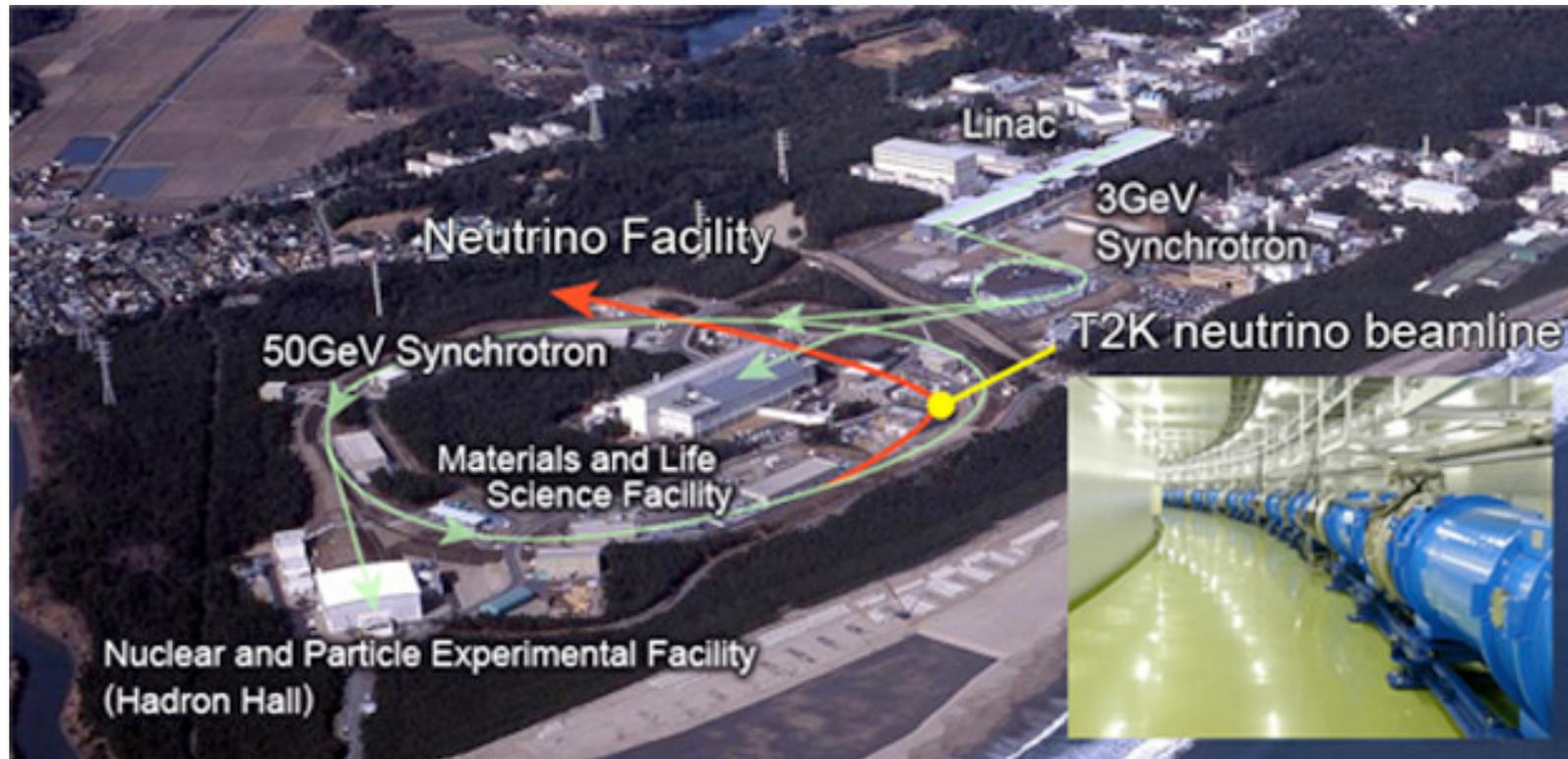
## CP violation, $\theta_{13}$ , leptogenesis

From modeling the early universe

- baryon number violation
- out of equilibrium interactions
- C and CP violation: known SM ~~CP~~ sources appear insufficient

but  $J_{CP}^\nu = \sin \theta_{12} \sin \theta_{23} \sin \theta_{13} \cos \theta_{12} \cos \theta_{23} \cos \theta_{13}^2 \sin \delta$   
 $\sim 0.2 \sin \theta_{13} \sin \delta$

T2K  $\nu_\mu \rightarrow \nu_e$  search



## Experimental parameters

- $2.5^\circ$  off-axis relatively narrow  $\nu$  beam, yielding  $E_\nu^{\text{peak}} \sim 0.6 \text{ GeV}$
- the J-PARC : SuperK baseline, which then places the detector at the  $\Delta m_{23}$  first oscillation maximum
- $\nu_\mu \rightarrow \nu_e$  appearance at a baseline much shorter than that optimizing appearance via  $\theta_{12}$ , so the effects of  $\theta_{13}$  can be seen

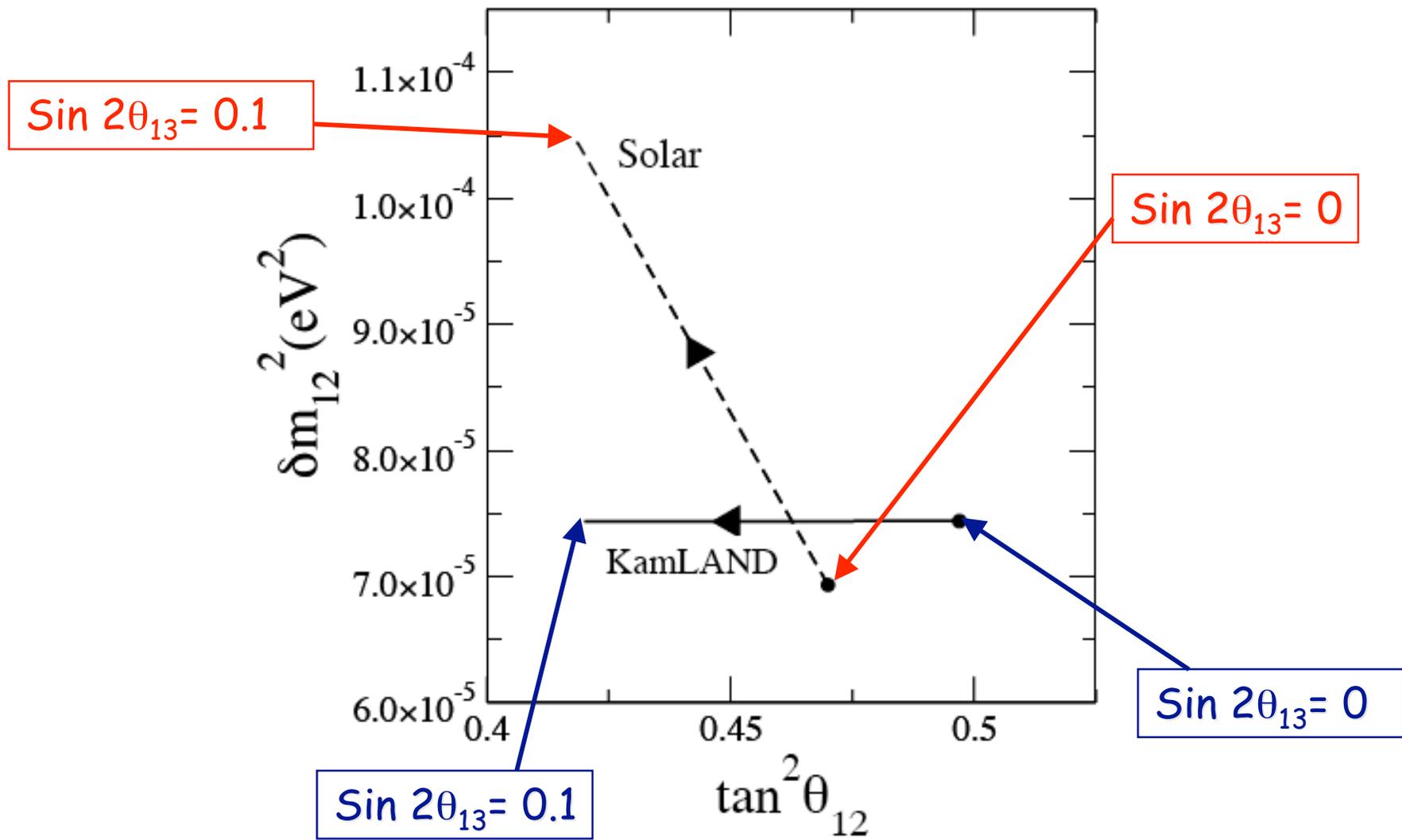
find 6 events when  $1.5 \pm 0.3$  would be expected were  $\theta_{12} = 0$  ( $2.5 \sigma$ )

deduce  $0.03(0.04) \lesssim \sin^2 2\theta_{13} \lesssim 0.28(0.34)$  normal(inverted),  $\delta_{\text{CP}} = 0$

best value, normal hierarchy  $\sim 0.11$

compares to CHOOZ, MINOS  $\sin^2 2\theta_{13} \lesssim 0.15$  and potentially indicates significant future LBNE sensitivity to  $\delta_{\text{CP}} = 0$

Probably need to wait: values this large may be problematic elsewhere



## LBNE: hierarchy and CP violation

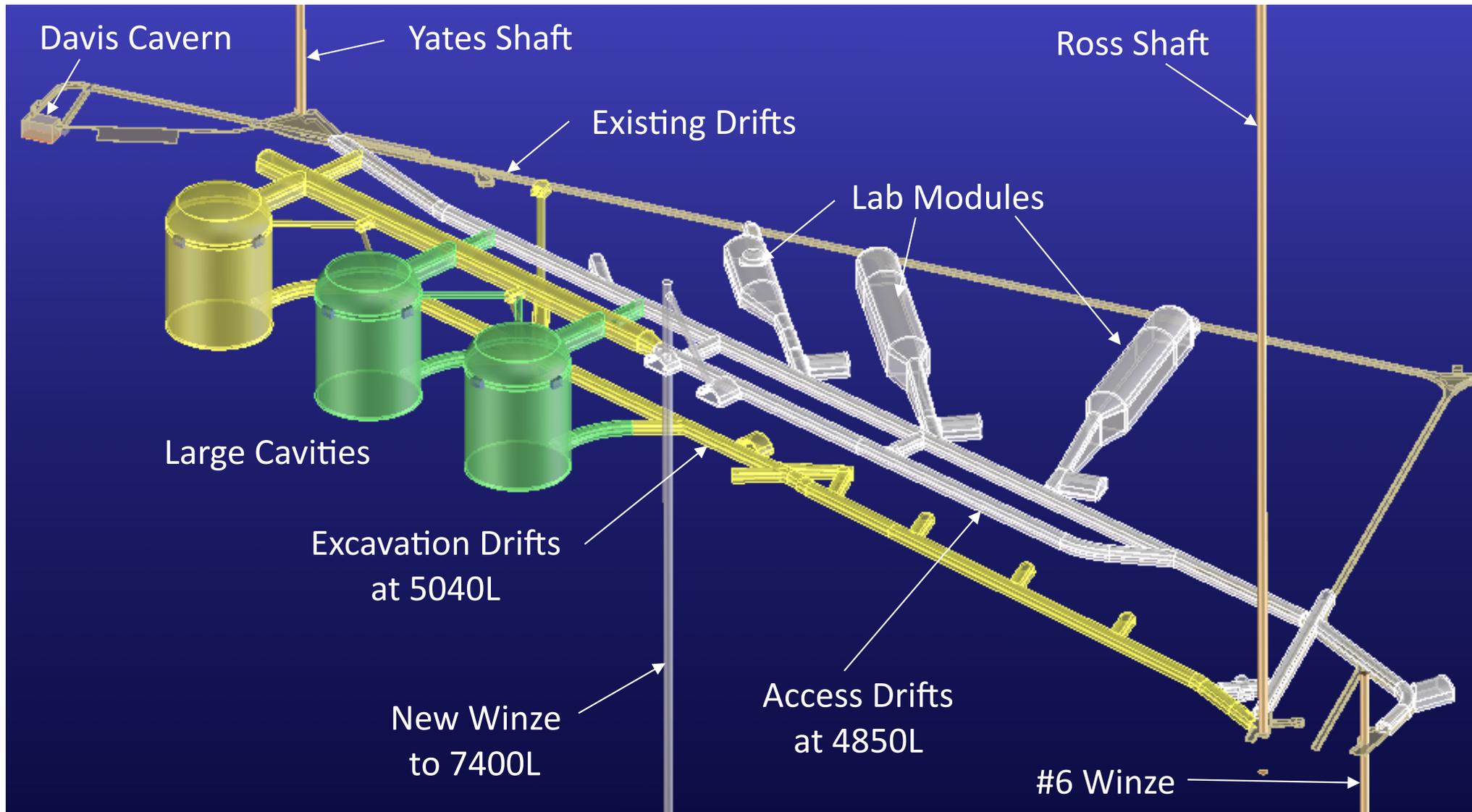


700 kW beam, on axis, water/argon megadetector, beamline to “DUSEL”

1300 km of matter: sign of matter effects  $\Leftrightarrow$  normal/inverted;

5 years of  $\nu_{\mu}S, \bar{\nu}_{\mu}S$  running  $\nu_{\mu} \rightarrow \nu_e$  vs  $\bar{\nu}_{\mu} \rightarrow \bar{\nu}_e$  for ~~CP~~

sounds like a good plan...



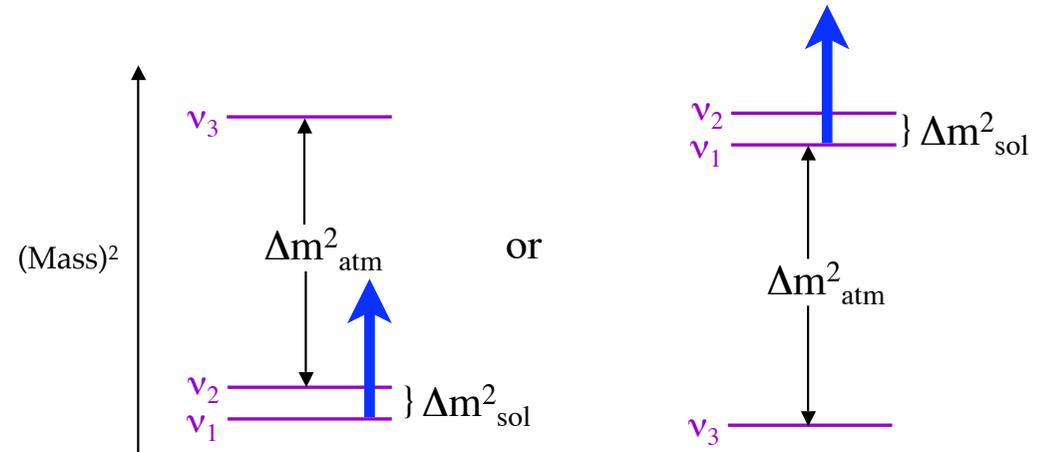
# Vacuum formula

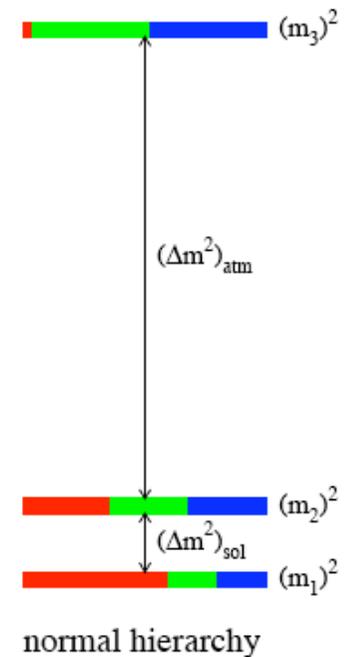
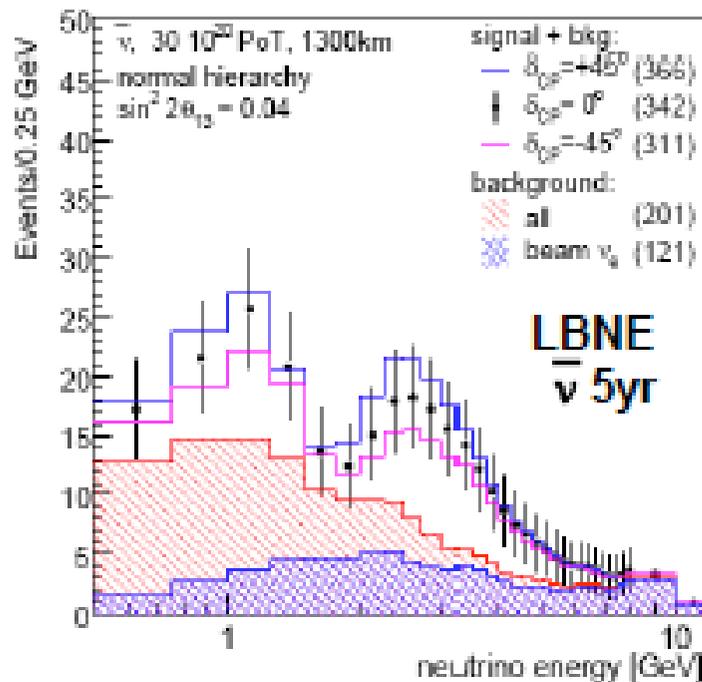
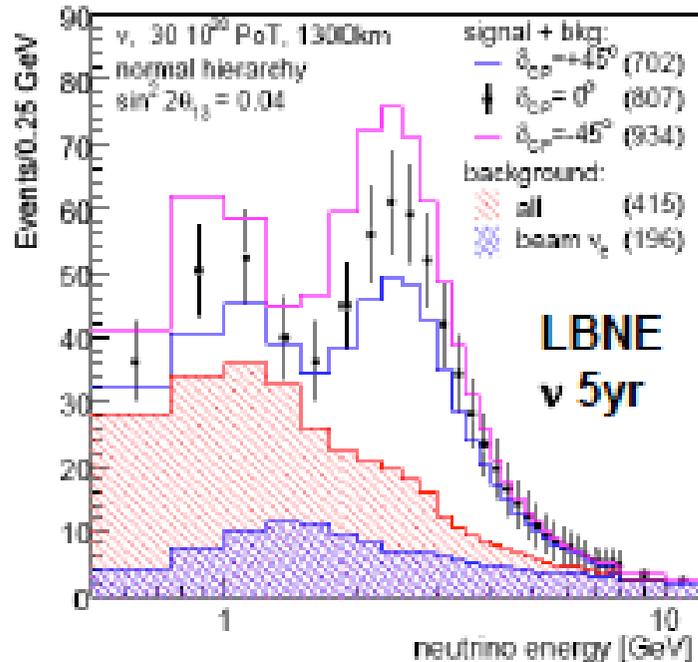
$$P \begin{pmatrix} \nu_\mu \rightarrow \nu_e \\ \bar{\nu}_\mu \rightarrow \bar{\nu}_e \end{pmatrix} = \frac{(\sin^2 2\theta_{23} \sin^2 2\theta_{13})(\sin^2 \Delta_{31})}{\pm \sin \delta (\sin 2\theta_{13} \sin 2\theta_{23} \sin 2\theta_{12})(\sin^2 \Delta_{31} \sin \Delta_{21})} + \frac{\cos \delta (\sin 2\theta_{13} \sin 2\theta_{23} \sin 2\theta_{12})(\sin \Delta_{31} \cos \Delta_{31} \sin \Delta_{21})}{+(\cos^2 \theta_{23} \sin^2 2\theta_{12})(\sin^2 \Delta_{21})}$$

nonzero?

altered by matter

Effects intertwined, as two channels are not CP conjugate when in matter



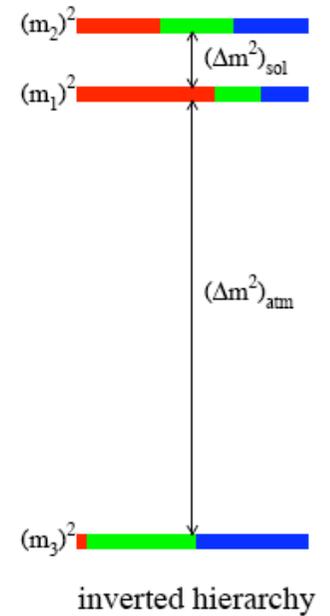
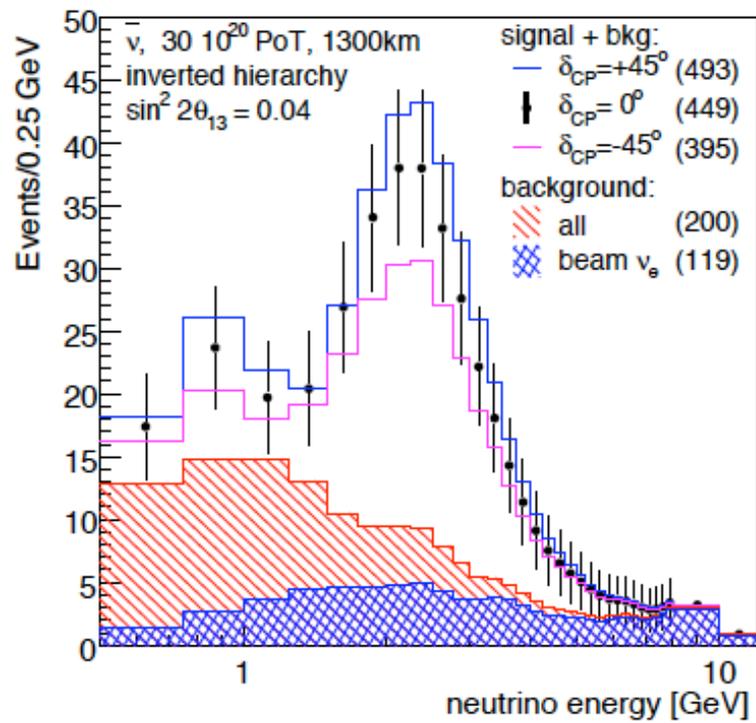
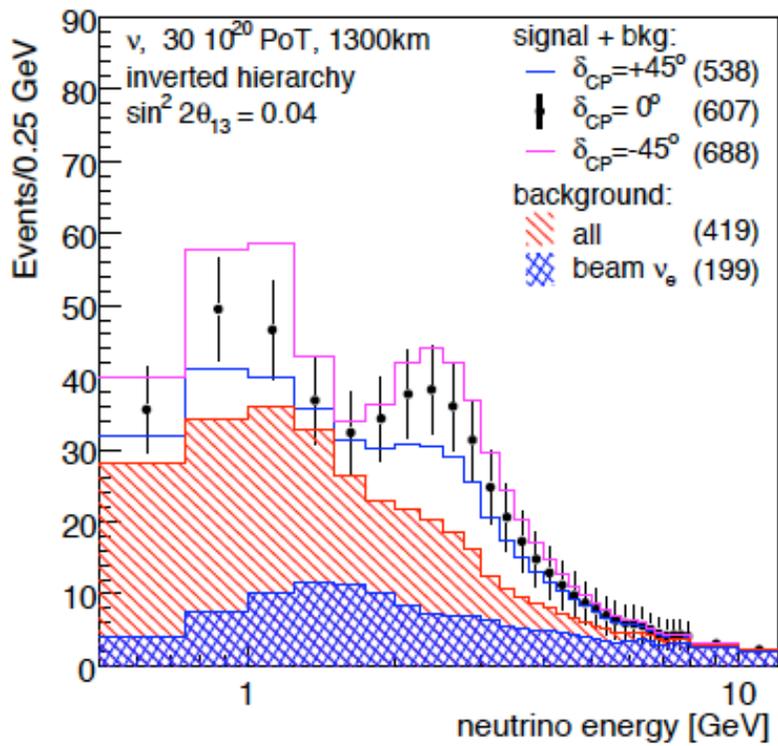


challenging:

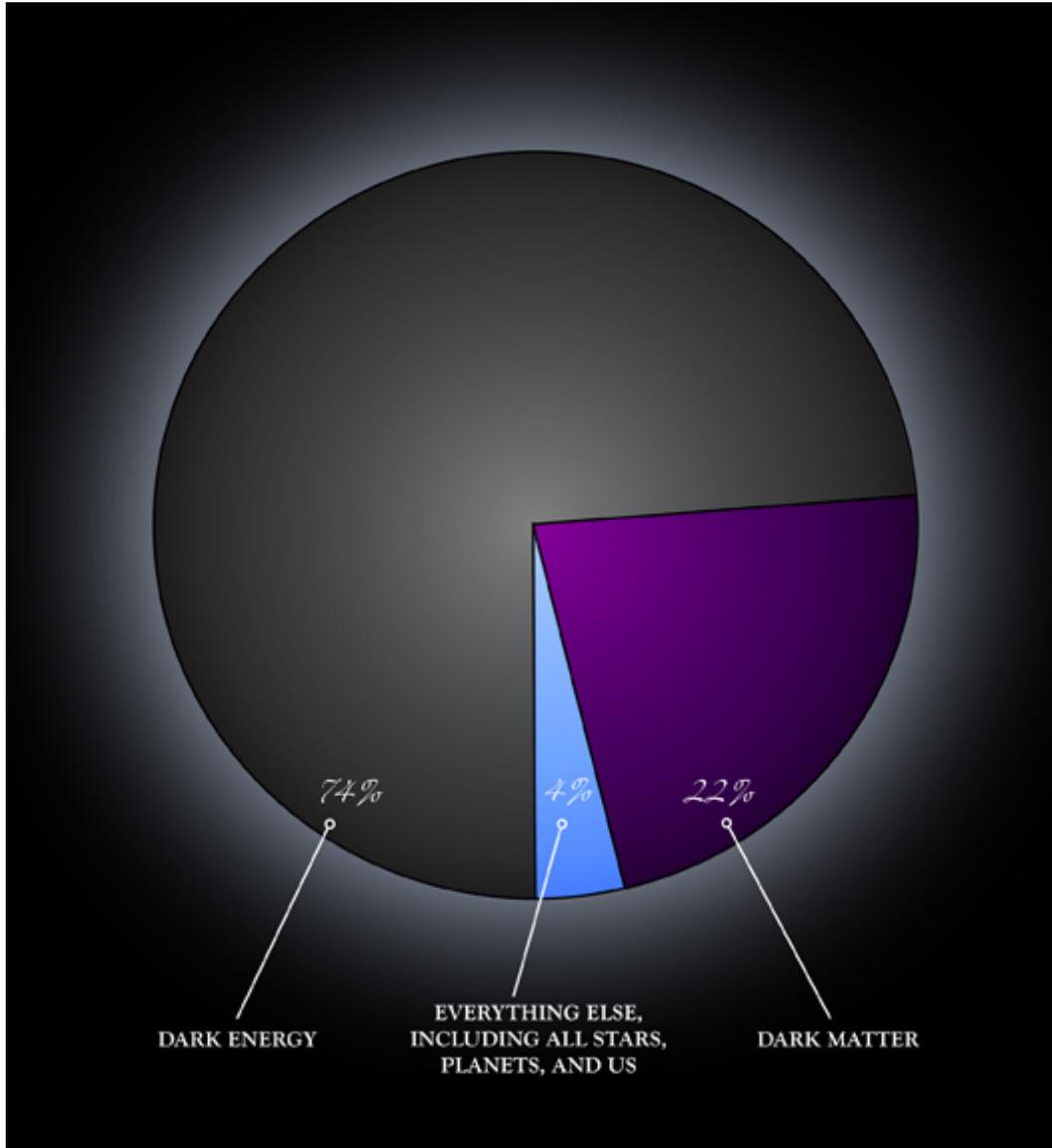
broad band beam; baseline requires a spectrum centered at about 2 GeV

low statistics, some beam contamination, backgrounds from  $\pi^0$  production

must be able to identify events (quasielastic kinematics) for which one can reconstruct the initial beam energy



## The Inner Space ↔ Outer Space Connection



To date, the major discoveries have come from astro/cosmo

uncertain properties of SM vs known/suspected to influence

- baryon asymmetry
- BBN: the number of relativistic degrees of freedom, net lepton number
- DM density, DM effects on expansion
- unique astrophysical oscillation environments
  - novel MSW effects
  - new level crossings

Plus incomparable sensitivity to new (sterile) vs

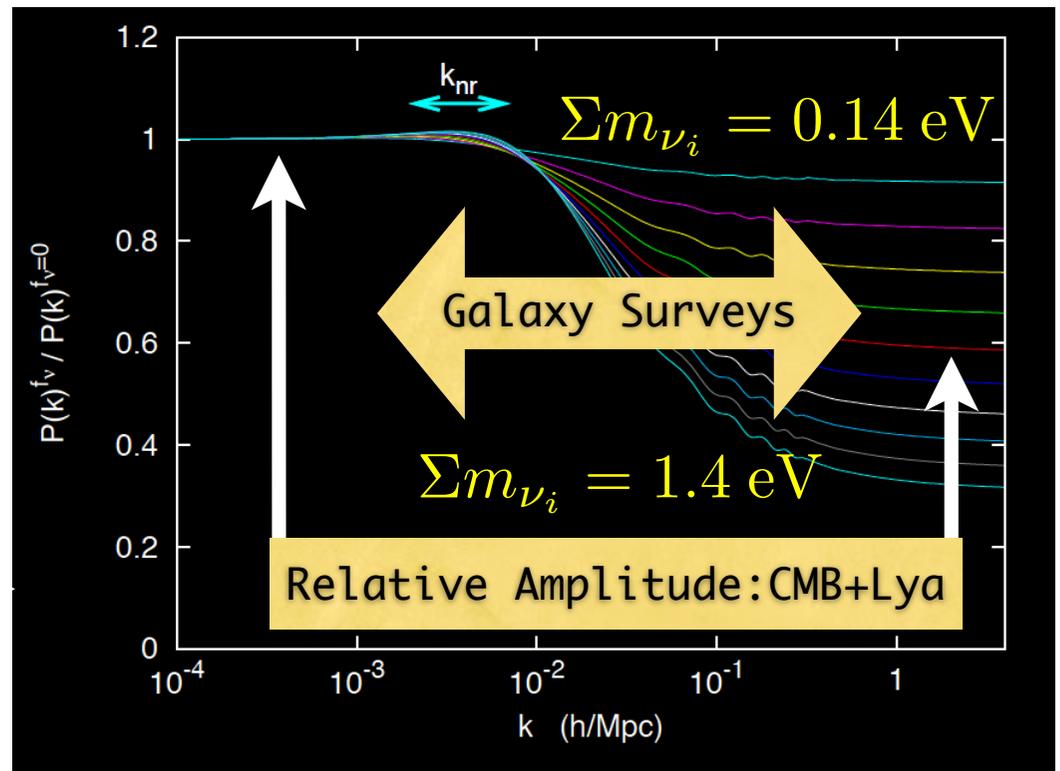
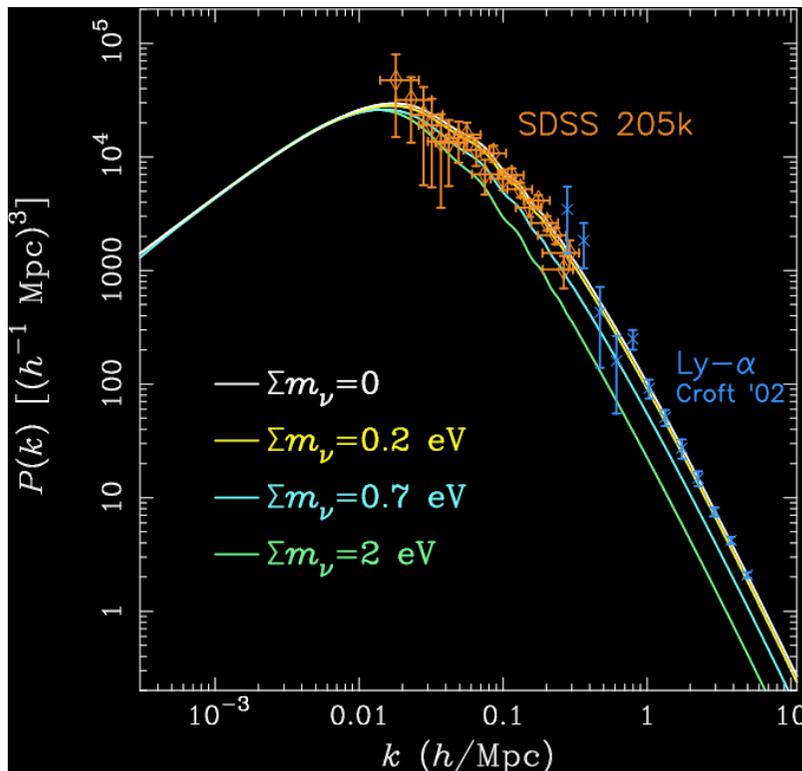
Absolute  $\nu$  mass scale: the one “known” component of DM is the  $\nu$

Standard BBN:  $n_\nu = N_\nu \left( \frac{3}{11} \right) n_\gamma \sim 340/\text{cm}^3$

$$\rho_\nu = \sum m_i n_{\nu_i} \quad \Omega_\nu \sim \sum \frac{m_{\nu_i}}{h^2 93 \text{ eV}} \quad E_i^2 = p^2 + m_i^2$$

smaller mass  $\nu$ s  $\Rightarrow$  relativistic longer, travel further  $\Rightarrow$  suppress growth of structure on larger scales

$$k_{\text{free streaming}} \sim 0.004 \sqrt{m_\nu / 0.05 \text{ eV}} \text{ Mpc}^{-1}$$



from Kev Abazajian

Thus  $\nu$  influences on structure evolve with both redshift  $Z$  and spatial scale in a characteristic way:

$$\sum m_\nu \sim 0.05 \text{ eV}, \quad z = \begin{pmatrix} 3.5 \\ 3.5 \\ 1.5 \\ 0.0 \end{pmatrix} \Rightarrow \text{power decrease} \sim \begin{pmatrix} 1.9\% \\ 1.0\% \\ 2.1\% \\ 3.5\% \end{pmatrix} \text{ for } k > \begin{pmatrix} 0.6 \\ 0.03 \\ 0.6 \\ 0.6 \end{pmatrix} \frac{1}{\text{Mpc}}$$

leverage: alter baryons + CDM at the  $\sim$  % level, when  $\Omega_\nu \sim 0.1\%$

typical combined analysis using existing data

$$\sum m_{\nu_i} < 0.58 \text{ eV} \quad \text{Komatsu et al. 2010, WMAP7 + SDSS LRG BAO + Ho}$$

$$\left( \frac{\Delta P}{P} \right)_{\text{future}} \sim 1\% \sim -12 \frac{\Omega_\nu}{\Omega_m} \Rightarrow \sum m_{\nu_i} \sim 11 \text{ meV}$$

Hu, Eisenstein, & Tegmark 1998; Abazajian & Dodelson 2003

A series of Astro2010 white papers were submitted that examined consequences of anticipated surveys, typically  $\sim \times 100$  increase in statistics

- high redshift galaxy surveys, SDSS-III BOSS  $10^5$  QSO survey, Planck CMB data, 21-cm radio telescopes with  $0.1 \text{ km}^2$  collection, weak lensing
- the statistical power for discovery at 50 meV were variously estimated at  $1-7\sigma$ , depending on the assumptions made on combining data sets

may be the **field's only nearterm** strategy for determining  $\sum m_\nu(i)$  for many scenarios, e.g., normal hierarchy with  $m_\nu(1) \sim 0$

this is also an example of a scenario where one could be sensitive to the **hierarchy**: inverted hierarchy requires  $\sum m_\nu(i) \gtrsim 10 \text{ meV}$

Systematic contributions to the error? Harder to assess

The leverage one gains from combining different types of measurements to increase the range in scale and  $Z$ , has the downside of increasing the chances of systematic conflicts

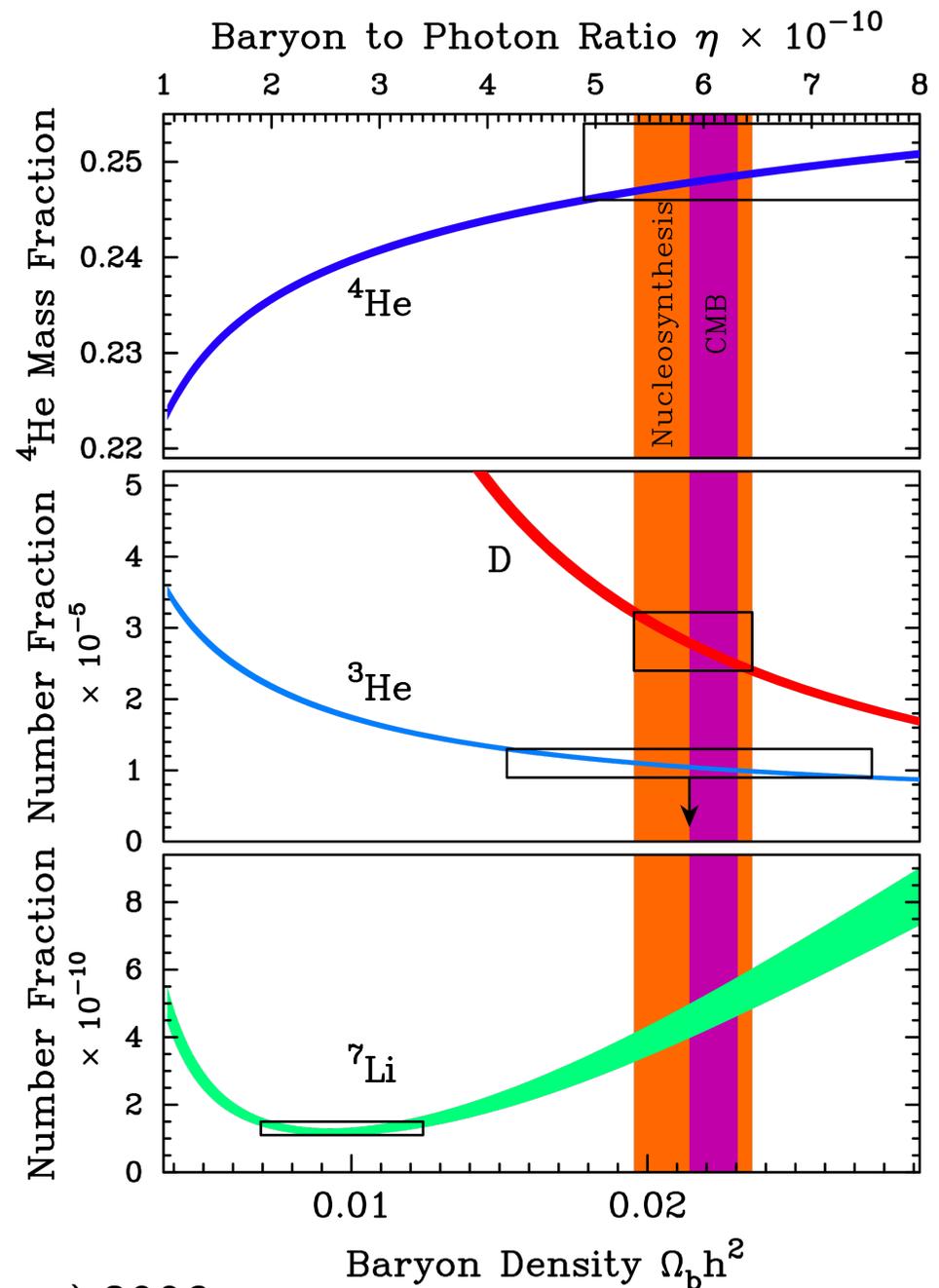
One of the reasons that the kind of unified computational program discussed here could be helpful: a team with the capacity to build, then continue to develop, a standard cosmological model, to fully vary that model, to apply it in a consistent way to disparate data sets, then to come to consensus when discrepancies among data sets emerge  
(reminiscent of the  $\Lambda$ CDM, with its 19 free parameters)

$m_\nu$  could be an inner space/outer space “home run,” impacting lab  $\nu$  physics interpretations:  $\beta\beta$  decay and LBNE (hierarchy)

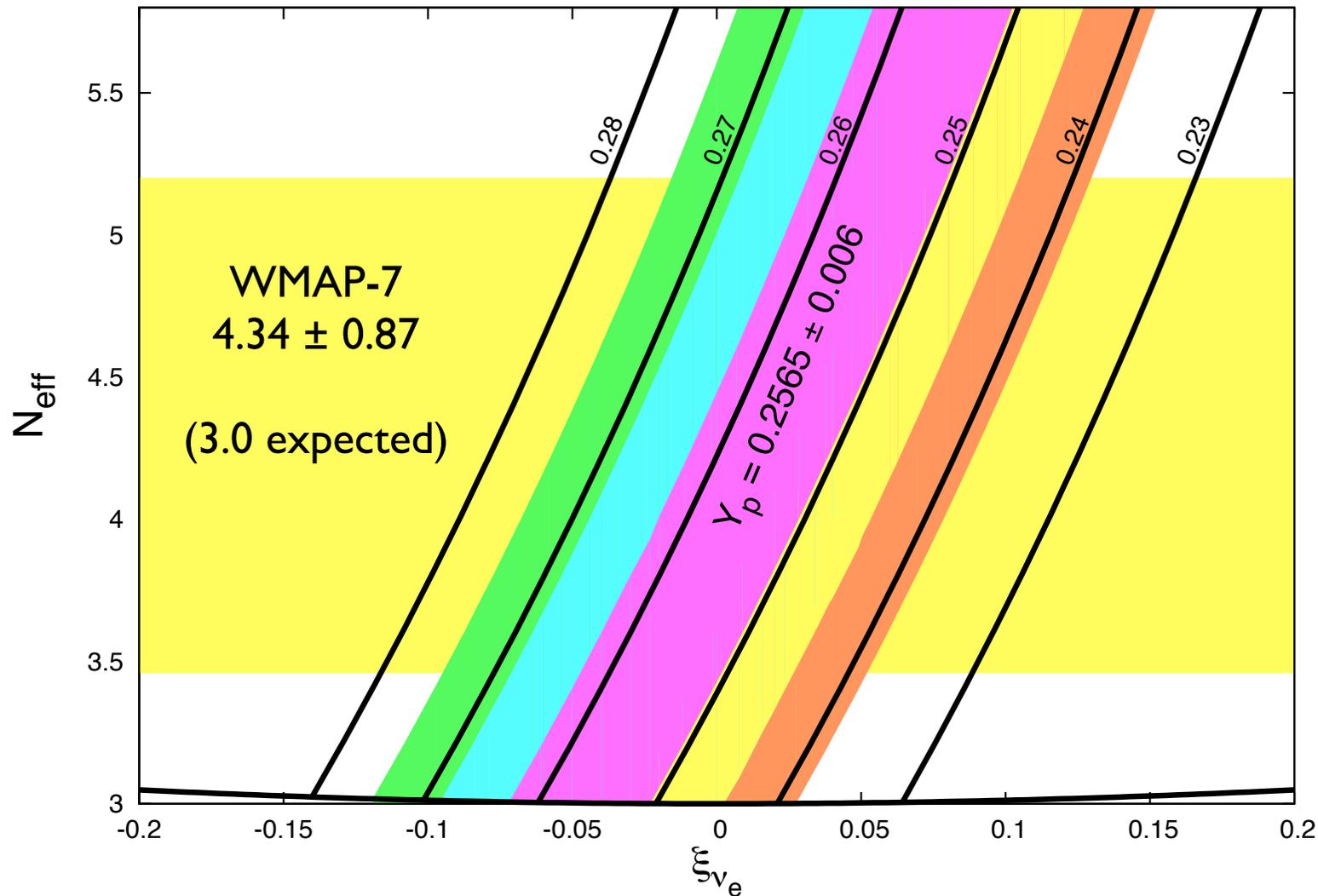
BBN: issues include  $\eta = n_B/n_\gamma$  consistency, the number of relativistic species (e.g., sterile neutrinos), the lepton number asymmetry, and alternatives to conventional abundance determinations

An issue exists with  ${}^7\text{Li}$ , which has a well-defined primordial abundance plateau, corresponding to an  $\eta \sim \eta_{\text{CMB}}$

The tension is  ${}^7\text{Li} - d$ , with cosmology indicating that  ${}^7\text{Li}$  is the outlier



# Competing clocks of expansion driven by the relativistic species and weak interactions driving $n$ densities downward



BBN and CMB studies constrain the  $\nu$  number and asymmetry  
weak hints that all is not right (but best to wait for Planck...)

Abundances: Potential to cross-check conventional low-Z abundance determinations with high-Z, pre-stellar determinations

$^4\text{He}$  at recombination provides an earlier sink for electrons, altering the electron density and thus the radiation scattering, and thus spectrum

(and  $^7\text{Li}$  provides a tiny amount of reheating)

$$Y_{\text{P}} = 0.326 \pm 0.075 \quad (\text{Komatsu et al. 2010: WMAP7 + BAO + Ho})$$

Deuterium abundance deduced from QSO absorption line systems  $\Rightarrow$

$$\eta_{\text{BBN}} \Rightarrow \text{assuming } N_{\nu} = 3 \Rightarrow Y_{\text{P}} = 0.2482 \pm 0.0007$$

Direct stellar determinations  $0.25 \pm 0.004$

## Comments

The analogy between the proposal to better organize computational cosmology, in response to rapidly advancing observations, is reminiscent to the Fowler/Bahcall/Iben/Sears initiative in 1962, to develop the standard solar model in anticipating of future neutrino experiments

Provided a very important way to correlate data, to feed in steadily improving microphysics (opacities, nuclear cross sections), to improve the physics when new measurements (e.g. helioseismology) required this

Led to major discoveries

The scales of physics are different today, the computers more complex, the computational teams much larger ... but the essential role of modeling in data-driven fields remains